# **1.2.4 Polarization-Mode Dispersion**

- Contraction of the second
- A single mode fiber is not a truly single mode fiber because it can support two degenerate modes that are that are polarized in two orthogonal directions
- Under ideal conditions (perfect cylindrical symmetry and a stress-free fiber),
  - a mode excited with its polarization in the *x* direction would not couple to the mode with the orthogonal y-polarization state.
- In real fibers, small departures from cylindrical symmetry, occurring because of random variations in the core shape along the fiber length,
  - result in a mixing of the two polarization states by breaking the mode degeneracy.
- > The **stress-induced anisotropy** can also break this degeneracy.

# Polarization in a single mode fiber



#### https://www.youtube.com/watch?v=kuht5Nv3Iio







> Modal birefringence: the mode-propagation constant  $\beta$  becomes slightly different for the modes polarized in the x and y directions.

The strength of modal birefringence is defined by

$$B_m = \frac{\left|\beta_x - \beta_y\right|}{k_0} = \left|n_x - n_y\right|$$

For a given value of  $B_m$ , the **two modes** exchange their powers in a **periodic fashion** as they propagate inside the fiber with the period.

$$L_{B} = \frac{2\pi}{\left|\beta_{x} - \beta_{y}\right|} = \frac{\lambda}{B_{m}}$$

 $L_{\rm B}$ : the beat length

**Fast axis:** the axis along which the **mode index is smaller** because the group velocity is larger for light propagating in that direction. **Slow axis:** the axis with the **larger mode index** 

- > In standard optical fibers,  $B_m$  is not constant along the fiber but changes randomly because of fluctuations in the core shape and anisotropic stress.
  - light launched into the fiber with a **fixed state of polarization** changes its polarization in a **random fashion**.
- This change in polarization is typically harmless for continuouswave (CW) light because most photodetectors do not respond to polarization changes of the incident light.
- It becomes an issue for optical communication systems when short pulses are transmitted over long lengths.
  - If an input pulse excites both polarization components, the **two components** travel along the fiber **at different speeds** because of their **different group velocities**.

## **General Description**



- These manual polarization controllers utilize stress-induced birefringence to alter the polarization in single mode fiber that is looped around two or three independent spools to create two or three independent fractional wave plates (fiber retarders).
- To transform an arbitrary input polarization state into an arbitrary output polarization state, a combination of
  - three paddles (a quarter-wave plate, a halfwave plate, and a quarter-wave plate) or
  - two paddles (quarter-wave plate and a quarter-wave plate) is used.



## **Polarization Controller**



#### https://www.youtube.com/watch?v=I9kvcN1Ew1k





The retardance of each paddle may be estimated from the following equation:

$$\varphi(Radians) = \frac{2\pi^2 a N d^2}{\lambda D}$$

- $\phi$  is the retardance,
- -a is a constant (0.133 for silica fiber),
- N is the number of loops,
- *d* is the fiber cladding diameter,
- $-\lambda$  is the wavelength, and
- *D* is the loop diameter

#### **3-Paddle Fiber Polarization Controllers**



Item #	FPC030	FPC031	FPC032	
Paddle Material	Black Delrin			
Number of Paddles		3		
Loop Diameter		1.06" (27 mm)	)	
Paddle Rotation		±117.5°		
Foot Print (L x W)	8.5" x 1.0" (215.9 mm x 25.4 mm)			
Fiber	None	CCC1310-J9		
Operating Wavelength Range <sup>a</sup>	N/A	1260 - 1625 nm		
Design Wavelength <sup>b</sup>	N/A	1310 nm		
Mode Field Diameter	N/A	8.6 ± 0.4 μm @ 1310 nm 9.7 ± 0.5 μm @ 1550 nm		
Cladding Diameter	N/A	125 ± 0.7 μm		
Coating Diameter	N/A	242 ± 5 μm		
Tubing Diameter	N/A	Ø900 µm Tight Buffer		
Numerical Aperture	N/A	0.14		
Loop Configuration <sup>c</sup>	N/A 2-3-2			
Connectors	N/A	FC/PC FC/APC		
Bend Loss	N/A	≤0.1 dB		

 Retardance varies as a function of wavelength. Refer to Chapter 2 for more information.

b. Devices with preloaded fiber are optimized for this wavelength.

c. For polarization controllers with fiber preinstalled.

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ltem #	FPC560	FPC561	FPC562	
Paddle Material	Black Delrin			
Number of Paddles		3		
Loop Diameter	2.2" (56 mm)			
Paddle Rotation	±117.5°			
Foot Print (L x W )	12.5" x 1.0" (317.5 mm x 25.4 mm)			
Fiber	None	SMF-28-J9		
Operating Wavelength Range <sup>a</sup>	N/A	1260 - 1625 nm		
Design Wavelength <sup>b</sup>	N/A	1310 nm		
Mode Field Diameter	N/A 9.2 ± 0.4 µm @ 1310 nm			
		10.4 ± 0.5 μm	@ 1550 nm	
Cladding Diameter	N/A	125 ± 0.7 μm		
Coating Diameter	N/A	242 ± 5 μm		
Tubing Diameter	N/A	Ø900 µm Tight Buffer		
Numerical Aperture	N/A	0.14		
Loop Configuration <sup>c</sup>	N/A	3-6-3		
Connectors	N/A	FC/PC FC/PC		
Bend Loss	N/A	≤0.1 dB		

a. Retardance varies as a function of wavelength. Refer to Chapter 2 for more information.

b. Devices with preloaded fiber are optimized for this wavelength.

c. For polarization controllers with fiber preinstalled.





The retardation of multi-order (including zero order) quarterwave plate is given by the following equation:

where m is an integer.

 $(2m + 1)\pi$ 

The retardation of multi-order (including zero order) halfwave plate is given by:

 $(2m+1)\pi$ 

The retardation of each paddle should be close to any number.

Order	m	Quarter-Wave Plate Retardation	Half-Wave Plate Retardation
Zero	0	$\frac{\pi}{2} \approx 1.57$	$\pi pprox 3.14$
1 <sup>st</sup>	1	$\frac{3\pi}{2} \approx 4.71$	$3\pi \approx 9.42$
2 <sup>nd</sup>	2	$\frac{5\pi}{2} \approx 7.85$	$5\pipprox$ 15.71
3 <sup>rd</sup>	3	$\frac{7\pi}{2} \approx 11.00$	$7\pipprox$ 21.99
4 <sup>th</sup>	4	$\frac{9\pi}{2} \approx 14.14$	$9\pipprox28.27$
5 <sup>th</sup>	5	$\frac{11\pi}{2} \approx 17.28$	$11\pi \approx 35.56$

The number of recommended loops and recommended fiber for several wavelengths is given in the following tables. These combinations come close to the desired quarter-wave retardation:

	# of Loops for ~1/4 $\lambda$ Retardation			
Wavelength	Ø18 mm	Ø27 mm	Ø56 mm	Recommended Fiber
480 nm	3 loops	N/A	3 loops	460HP, S450, S460
630 nm	3 loops	2 loops	4 loops	630HP or S630
850 nm	3 loops	6 loops	2 loops	780HP, SM800-5.6
980 nm	2 loops	3 loops	_2 loops	_980HP_HI1060_19_HI980_19
1060 nm	2 loops	3 loops	2 loops	980HP, HI1060-J9, HI980-J9
1310 nm	3 loops	2 loops	3 loops	SMF28e+ and CCC1310-J9

**FPC030** 

These combinations come close to the desired half-wave retardation:

	# of Loops for ~1/2 $\lambda$ Retardation			
Wavelength	Ø18 mm	Ø27 mm	Ø56 mm	Recommended Fiber
480 nm	2 loops	3 loops	2 loops	460HP, S450, S460
630 nm	1 loop	4 loops	3 loops	630HP or S630
850 nm	1 loop	2 loops	4 loops	780HP, SM800-5.6
980 nm	4 loops	2 loops	4 loops	980HP_HI1060-J9_HI980-J9_
1060 nm	3 loops	2 loops	5 loops	980HP, HI1060-J9, HI980-J9
1310 nm	2 loops	3 loops	6 loops	SMF28e+ and CCC1310-J9

#### **FPC560**

## **Electronic Polarization Controller**





Phoenix Photonics Electronic Polarization Controller

BIRD

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0.48 / 2.49

# separate control of knobs will give you full coverage of the Poincare shpere

Electronic Polarization Controller



- The pulse becomes broader at the output end because group velocities change randomly in response to random changes in fiber birefringence (analogous to a random-walk problem).
- This phenomenon, referred to as polarization-mode dispersion (PMD 偏振模色散), has been studied extensively because of its importance for long-haul lightwave systems.

## **Polarization Mode Dispersion**





#### **From YouTube**

- > The extent of pulse broadening can be estimated from the time delay  $\Delta T$  occurring between the two polarization components during propagation of an optical pulse.
- $\succ \Delta T$  is given by:



- $\Delta\beta_1$ : related to group-velocity mismatch
- *L* : fiber of length
- $B_{\rm m}$ : constant birefringence.
- Equation (1.2.16) cannot be used directly to estimate PMD for standard telecommunication fibers because of random changes in birefringence occurring along the fiber.
- These changes tend to equalize the propagation times for the two polarization components.

- ➢ PMD is characterized by the root-mean-square (RMS 均方根) value of ∆T obtained after averaging over random perturbations.
- > The variance of  $\Delta T$  is found to be.

$$\sigma_T^2 = \left\langle (\Delta T)^2 \right\rangle = 2(\Delta \beta_1 l_c)^2 [\exp(-L/l_c) + L/l_c - 1]$$

- $\Delta\beta_1 \equiv \Delta\tau/L$ ,
- $\Delta \tau$ : the **differential group delay** along the **principal states** of polarization,
- *l<sub>c</sub>* (correlation length): the length over which two polarization components remain correlated;

For L > 0.1 km, we can use  $l_c$  (typical order of 10 m) <<L to find that

$$\sigma_T \approx \Delta \beta_1 \sqrt{2l_c L} \equiv D_p \sqrt{L}$$

 $D_{\rm p}$  : PMD parameter

## Polarization-maintaining or polarization preserving fibers:



- It is desirable that fibers transmit light without changing its state of polarization for some applications
- A large amount of birefringence is introduced intentionally in these fibers through design modifications so that relatively small birefringence fluctuations are masked by it and do not affect the state of polarization significantly.
- One scheme breaks the cylindrical symmetry by making the fiber core elliptical in shape.
  (The degree of birefringence is typically ~10<sup>-6</sup>.)
- An alternative scheme makes use of stress-induced birefringence (permits  $B_m \sim 10^{-4}$ .)

## **Schematic setup of PM fiber**



https://www.thorlabs.com/newgrouppage9.cfm?objectgroup\_id=159 6&pn=PS-PM980#7707 Thorlabs website Link

### **PM fiber**

![](_page_19_Picture_1.jpeg)

#### https://www.youtube.com/watch?v=7rrb-\_Iin-g

![](_page_19_Figure_3.jpeg)

![](_page_20_Picture_0.jpeg)

![](_page_20_Picture_1.jpeg)

### Specifications

Geometrical & Mechanical		
Core Diameter	5.5 µm	
Cladding Diameter	125 ± 2 μm	
Coating Diameter	245 ± 15 μm	
Core-Clad Offset	<0.5 µm	
Coating Concentricity	≤5 μm	
Coating Material	UV Cured, Dual Acrylate	
Operating Temperature	-40 to 85 °C	
Proof Test Level	200 kpsi (1.4 GN/m <sup>2</sup> )	

Optical		
Numerical Aperture	0.12	
Attenuation	≤2.5 dB/km @ 980 nm	
Operating Wavelength	970 - 1550 nm	
Second Mode Cut-off	920 ± 50 nm	
Mode Field Diameter (1/e <sup>2</sup> fit - near field)	6.6± 0.7 μm @ 980 nm	
Beat Length	≤2.7 mm @ 980 nm	
Normalized Cross Talk	≤-40 dB @ 4 m	
	≤-30 dB @ 100 m (nominal)	

- The use of polarization-maintaining fibers (保偏光纖) requires identification of the slow and fast axes before an optical signal can be launched into the fiber.
  - Structural changes are often made to the fiber for this purpose.
  - D fiber: cladding is flattened in such a way that the flat surface is parallel to the slow axis of the fiber.

![](_page_21_Figure_3.jpeg)

**Cladding-pumped fiber amplifier** based on a **double-clad fiber**. The signal light is launched into the doped core, while the pump light is launched into the inner cladding. The core is **D-shaped** for more efficient pump absorption. When the polarization direction of the linearly polarized light coincides with the slow or the fast axis, the state of polarization remains unchanged during propagation.

![](_page_22_Figure_1.jpeg)

Figure 1.8: Variation of birefringence parameter  $B_m$  with thickness d of the stress-inducing element for four different polarization-maintaining fibers. Different shapes of the stress-applying elements (shaded region) are shown in the

inset.

- If the polarization direction makes an angle with these axes, polarization changes continuously along the fiber in a periodic manner with a period equal to the beat length [see Eq. (1.2.15)].
- The state of polarization changes over one-half of the beat length from linear to elliptic, elliptic to circular, circular to elliptic, and then back to linear but is rotated by 90° from the incident linear polarization.

![](_page_23_Figure_2.jpeg)

Evolution of state of polarization along a polarization-maintaining fiber when <sup>24</sup> input signal is linearly polarized at 45 ° *from the slow axis*.

#### Wavelength-Tunable High-Energy All-Normal-Dispersion Yb-Doped Mode-Locked All-Fiber Laser with a HiBi Fiber Sagnac Loop Filter

![](_page_24_Picture_1.jpeg)

- The resulting transmission spectrum of the HiBi fiber Sagnac loop is approximately a periodic function of the wavelength given as
  - $T = (1/2)[1 \cos(2\pi BL/\lambda)]$
- L, B, and λ are the length of the HiBi fiber the HiBi fiber birefringence, and the wavelength in vacuum

![](_page_24_Figure_5.jpeg)

- ➤ The birefringence can be expressed as  $B = \lambda_0 / L_B,$   $L_B: \text{ beat length measured at wavelength } \lambda_0.$
- The peak spacing is given as  $\Delta \lambda = \lambda^2 / (BL)$

![](_page_25_Figure_0.jpeg)

Photonic Technology Lab.

## Lyot Filter in sensor [1]

![](_page_26_Figure_1.jpeg)

[1] Huang et al, "Intensity modulated torsion sensor based on optical fiber reflective Lyot filter Opt. Express, 25, 5081 (2017)

## Intensity modulated torsion sensor

![](_page_27_Picture_1.jpeg)

![](_page_27_Figure_2.jpeg)

Schematic diagram of the experiment setup for the reflective all-fiber Lyot filter-based torsion sensing system.

![](_page_27_Figure_4.jpeg)

Reflection spectra of the Lyot filter under torsion of  $10^{\circ}$ ,  $20^{\circ}$ ,  $30^{\circ}$ ,  $40^{\circ}$ ,  $50^{\circ}$ .

![](_page_27_Picture_6.jpeg)

# **1.3 Fiber Nonlinearities**

![](_page_28_Figure_1.jpeg)

- The response of any dielectric to light becomes nonlinear for intense electromagnetic fields, and optical fibers are no exception.
- The origin of nonlinear response is related to aharmonic motion of bound electrons under the influence of an applied field.
- The total polarization P induced by electric dipoles is not linear in the electric field E, but satisfies the more general relation.

$$P = \varepsilon_0(\chi^{(1)} \cdot E + \chi^{(2)} \vdots EE + \chi^{(3)} \vdots EEE + \cdots)$$

 $\varepsilon_0$ : the vacuum permittivity  $\chi^{(j)}$ :  $j_{th}$  order susceptibility (a tensor of rank j+1)